

Contexts in Convex and Sequential Effect Algebras

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A convex sequential effect algebra (COSEA) is an algebraic system with three physically motivated operations, an orthogonal sum, a scalar product and a sequential product. The elements of a COSEA correspond to yes-no measurements and are called effects. In this work we stress the importance of contexts in a COSEA. A context is a finest sharp measurement and an effect will act differently according to the underlying context with which it is measured. Under a change of context, the possible values of an effect do not change but the way these values are obtained may be different. In this paper we discuss direct sums and the center of a COSEA. We also consider conditional probabilities and the spectra of effects. Finally, we characterize COSEA's that are isomorphic to COSEA's of positive operators on a complex Hilbert space. These result in the traditional quantum formalism. All of this work depends heavily on the concept of a context.

1 Introduction

We present an axiomatic framework for quantum mechanics in which the basic entities and operations have physical significance. In this framework, the principle concepts are states and effects. The states represent initial preparations that describe the condition of the system, while the effects represent yes-no measurements that probe the system. The effects may be unsharp or fuzzy [5, 6, 9, 18]. A state applied to an effect produces the probability that the effect gives a yes value when the system is in that state. The resulting mathematical structure is called a convex sequential effect algebra (COSEA) \mathcal{E} [10, 15, 11, 21, 22]. The three mathematical operations in \mathcal{E} are an orthogonal sum $a \oplus b$, a scalar product $\lambda a, \lambda \in [0, 1] \subseteq \mathbb{R}$ and a sequential product $a \circ b$. These operations have physical interpretations that we now discuss.

Although this framework is much more general, we can employ the model of an optical bench to visualize what is happening here. A beam of particles (photons, electrons, etc.) is emitted from a source and propagates through a channel on the bench until the beam arrives at a detector at the end of the channel. The particles are initially prepared in a certain state and the effects describe various filters that can be placed in the channel. The beam travels through one or more filters which interact with the beam and can change its properties in certain ways. The detector may count particles or measure different characteristics of the beam. The sum $a \oplus b$ is performed by first splitting the beam into two equal parts, which are directed toward the two filters placed in parallel after which both beams are reunited before being collected at the detector. The scalar product λa corresponds to an attenuation of filter a by the factor λ . This can be accomplished by placing a gray filter with a certain darkness in front of filter a . The gray filter blocks some of the particles but does not otherwise disturb the beam. The sequential product $a \circ b$ is performed by placing the filters in series so that a is first and b is second. In this way, filter a can interfere with the operation of filter b while b cannot interfere with the operation of a . We will find this useful for describing quantum interference.

In this work, an important role will be played by the context under which an effect is observed. A context is a finest sharp measurement and an effect will act differently according to the underlying context with which it is measured. For example, in the optical bench scenario, changing contexts may result from altering the detectors or varying the size, shape or location of the bench. Under a change of context, the possible values of an effect do not change but the way these values are obtained may be different. As far as contexts are concerned, there is a great difference between classical and quantum systems. We shall show that classical systems have exactly one context, while quantum systems have infinitely many.

In Section 2 we define the concepts of COSEA's and contexts. Section 3 discusses direct sums and the center of a COSEA. Section 4 considers conditional probabilities and spectra of effects. Finally, Section 5 characterizes COSEA's that are isomorphic to COSEA's of positive operators on a complex Hilbert space. Of course, these result in the traditional quantum formalism. There is some overlap of this paper and the work in [21, 22]. However, our stress on contexts provides a different approach.

2 Convex Sequential Effect Algebras

Let \mathcal{E} be the set of effects and \mathcal{S} the set of states for a physical system. The connection between \mathcal{E} and \mathcal{S} is given by a *probability function* $F: \mathcal{E} \times \mathcal{S} \rightarrow [0, 1] \subseteq \mathbb{R}$ where $F(a, s)$ is interpreted as the probability that effect a has a yes value when the system is in state s . An *effect-state space* is a triple $(\mathcal{E}, \mathcal{S}, F)$ where \mathcal{E} and \mathcal{S} are nonempty sets and $F: \mathcal{E} \times \mathcal{S} \rightarrow [0, 1]$ satisfies:

- (ES1) There exist elements $0, 1 \in \mathcal{E}$ such that $F(0, s) = 0$, $F(1, s) = 1$ for every $s \in \mathcal{S}$.
- (ES2) If $F(a, s) \leq F(b, s)$ for every $s \in \mathcal{S}$, then there exists a unique $c \in \mathcal{E}$ such that $F(a, s) + F(c, s) = F(b, s)$ for all $s \in \mathcal{S}$.
- (ES3) If $a \in \mathcal{E}$ and $\lambda \in [0, 1]$, then there exists an element $\lambda a \in \mathcal{E}$ such that $F(\lambda a, s) = \lambda F(a, s)$ for all $s \in \mathcal{S}$.

The elements $0, 1$ in (ES1) correspond to the null effect that never occurs and the unit effect that always occurs, respectively. It is shown in [10, 15] that if $F(a, s) + F(b, s) \leq 1$ for every $s \in \mathcal{S}$, then there exists a unique $c \in \mathcal{E}$ such that

$$F(c, s) = F(a, s) + F(b, s)$$

for all $s \in \mathcal{S}$. We then write $a \perp b$ and define $a \oplus b = c$. In this way, \oplus is a partial binary operation on \mathcal{E} .

The structure $(\mathcal{E}, 0, 1, \oplus)$ is called an *effect algebra* and satisfies the following axioms:

- (EA1) If $a \perp b$, then $b \perp a$ and $b \oplus a = a \oplus b$,
- (EA2) If $a \perp b$ and $(a \oplus b) \perp c$, then $b \perp c$, $a \perp (b \oplus c)$ and $a \oplus (b \oplus c) = (a \oplus b) \oplus c$,
- (EA3) For every $a \in \mathcal{E}$ there exists a unique $a' \in \mathcal{E}$ such that $a \perp a'$ and $a \oplus a' = 1$,
- (EA4) If $a \perp 1$, then $a = 0$.

We define $a \leq b$ if there is a $c \in \mathcal{E}$ such that $a \oplus c = b$. The element c is unique and we write $c = b \ominus a$. It can be shown that $(\mathcal{E}, 0, 1, \leq)$ is a bounded poset and $a \perp b$ if and only if $a \leq b'$ [5, 6]. Moreover, $a'' = a$ and $a \leq b$ implies $b' \leq a'$ for all $a, b \in \mathcal{E}$. If we incorporate the scalar product λa of (ES3) we obtain the following structure. An effect algebra \mathcal{E} is *convex* [10, 15, 11] if for every $a \in \mathcal{E}$ and $\lambda \in [0, 1] \subseteq \mathbb{R}$ there exists an element $\lambda a \in \mathcal{E}$ such that

(CO1) If $\alpha, \beta \in [0, 1]$ and $a \in \mathcal{E}$, then $\alpha(\beta a) = (\alpha\beta)a$.

(CO2) If $\alpha, \beta \in [0, 1]$ with $\alpha + \beta \leq 1$ and $a \in \mathcal{E}$, then $\alpha a \perp \beta a$ and $(\alpha + \beta)a = \alpha a \oplus \beta a$.

(CO3) If $a, b \in \mathcal{E}$ with $a \perp b$ and $\lambda \in [0, 1]$, then $\lambda a \perp \lambda b$ and $\lambda(a \oplus b) = \lambda a \oplus \lambda b$.

(CO4) If $a \in \mathcal{E}$, then $1a = a$.

We call an effect algebra an EA and a convex effect algebra a COEA, for short. In \mathcal{E} and \mathcal{F} are EA's, a map $\phi: \mathcal{E} \rightarrow \mathcal{F}$ is *additive* if $a \perp b$ implies that $\phi(a) \perp \phi(b)$ and

$$\phi(a \oplus b) = \phi(a) \oplus \phi(b)$$

An additive map ϕ that satisfies $\phi(1) = 1$ is called a *morphism*. A morphism $\phi: \mathcal{E} \rightarrow \mathcal{F}$ for which $\phi(a) \perp \phi(b)$ implies $a \perp b$ is a *monomorphism* and a surjective monomorphism is an *isomorphism*. If \mathcal{E} and \mathcal{F} are COEA's, a morphism $\phi: \mathcal{E} \rightarrow \mathcal{F}$ is *affine* if $\phi(\lambda a) = \lambda \phi(a)$ for all $\lambda \in [0, 1]$, $a \in \mathcal{E}$. If there exists an affine isomorphism $\phi: \mathcal{E} \rightarrow \mathcal{F}$ we say that \mathcal{E} and \mathcal{F} are COEA *isomorphic*.

The simplest example of a COEA is the unit interval $[0, 1] \subseteq \mathbb{R}$ with the usual addition (when $a + b \leq 1$) and scalar multiplication. A *state* on an EA \mathcal{E} is a morphism $\omega: \mathcal{E} \rightarrow [0, 1]$. Notice that in an effect-state space, the function $a \mapsto F(a, s)$ is a state on \mathcal{E} . We denote the set of states on \mathcal{E} by $\Omega(\mathcal{E})$. We say that $S \subseteq \Omega(\mathcal{E})$ is *order-determining* if $\omega(a) \leq \omega(b)$ for all $\omega \in S$ implies that $a \leq b$. It is shown in [15] that every state on a COEA is affine. It is also shown in [15] that an effect-state space is equivalent to a COEA with an order-determining set of states.

We now introduce the sequential product $a \circ b$ on a COEA. Because of the series order for $a \circ b$, a may interfere with the b measurement but b will never interfere with the a measurement. If $a \circ b = b \circ a$ we write $a | b$ and say that a and b *do not interfere*. We now present our general definition.

A *convex sequential effect algebra* (COSEA) [11] is an algebraic system $(\mathcal{E}, 0, 1, \oplus, \circ)$ where $(\mathcal{E}, 0, 1, \oplus)$ is a COEA and $\circ: \mathcal{E} \times \mathcal{E} \rightarrow \mathcal{E}$ is a binary operation satisfying:

(S1) $b \mapsto a \circ b$ is additive for all $a \in \mathcal{E}$,

(S2) $1 \circ a = a$ for all $a \in \mathcal{E}$,

(S3) If $a \circ b = 0$, then $a | b$,

(S4) If $a | b$, then $a | b'$ and $a \circ (b \circ c) = (a \circ b) \circ c$ for all $c \in \mathcal{E}$,

(S5) If $c | a$ and $c | b$ then $c | a \circ b$ and $c | (a \oplus b)$ whenever $a \perp b$,

(S6) For all $\lambda \in [0, 1] \subseteq \mathbb{R}$, $a, b \in \mathcal{E}$, we have that

$$(\lambda a) \circ b = a \circ (\lambda b) = \lambda(a \circ b)$$

It is shown in [21] that if \mathcal{E} satisfies an additional continuity property that makes \mathcal{E} a σ -COSEA then (S6) is automatically satisfied.

In quantum mechanics, $a \circ b$ is useful for describing quantum interference. It is also needed for defining the important concept of conditional probability. An element a in a COSEA is *sharp* if the greatest lower bound $a \wedge a' = 0$. Sharp effects are thought of as effects that are precise or unfuzzy. We denote the set of sharp effects in \mathcal{E} by $S(\mathcal{E})$.

Theorem 2.1. [12] *The sequential product in a COSEA \mathcal{E} has the following properties. (i) $a \circ b \leq a$ for all $a, b \in \mathcal{E}$. (ii) If $a \leq b$, then $c \circ a \leq c \circ b$ for all $c \in \mathcal{E}$. (iii) $a \in S(\mathcal{E})$ if and only if $a \circ a = a$. (iv) For $a \in \mathcal{E}$, $b \in S(\mathcal{E})$, $a \circ b = 0$ if and only if $a \perp b$. (v) For $a \in \mathcal{E}$, $b \in S(\mathcal{E})$, $a \leq b$ if and only if $a \circ b = b \circ a = a$ and $b \leq a$ if and only if $a \circ b = b \circ a = b$.*

For a COSEA \mathcal{E} , we call $a \in S(\mathcal{E})$ *one-dimensional* if $a \neq 0$ and if $b \in \mathcal{E}$ with $b \leq a$, then $b = \lambda a$ for some $\lambda \in [0, 1]$. We denote the set of one-dimensional elements of \mathcal{E} by $S_1(\mathcal{E})$. It is shown in [11] that if $a \in S_1(\mathcal{E})$ then there exists an $\hat{a} \in \Omega(\mathcal{E})$ such that $\hat{a}(a) = 1$. A COSEA is *state-unique* if \hat{a} is unique. Although it is not known whether an arbitrary COEA is state-unique, it is shown in [21, 22] that every COSEA is state-unique.

A *finite context* in a COEA \mathcal{E} is a finite set $\{a_1, \dots, a_n\} \subseteq S_1(\mathcal{E})$ such that

$$a_1 \oplus a_2 \oplus \dots \oplus a_n = 1$$

It follows that $\hat{a}_i(a_j) = \delta_{ij}$. We denote the set of finite contexts in \mathcal{E} by $\mathcal{C}(\mathcal{E})$. We interpret a finite context as a finest sharp measurement. We say that \mathcal{E} is *finite-dimensional* if there does not exist an infinite sequence $a_i \in S_1(\mathcal{E})$ such that $a_1 \oplus \dots \oplus a_n$ is defined for all n . Thus, there are no infinite contexts. For simplicity, we assume that the COEA's (and COSEA's) we consider in this paper are finite-dimensional. If \mathcal{E} is state-unique and $a, b \in S_1(\mathcal{E})$, we call $\hat{a}(b)$ the *transition probability* from a to b . We say that \mathcal{E} is *symmetric* if $\hat{a}(b) = \hat{b}(a)$ for all $a, b \in S_1(\mathcal{E})$. It is shown in [21, 22] that every COSEA is symmetric.

Lemma 2.2. *If \mathcal{E} is state-unique and symmetric, then all contexts in \mathcal{E} have the same cardinality.*

Proof. Let $\mathcal{A}, \mathcal{B} \in \mathcal{C}(\mathcal{E})$ with $\mathcal{A} = \{a_1, \dots, a_n\}$, $\mathcal{B} = \{b_1, \dots, b_m\}$. Then

$$n = \sum_{i,j} \hat{a}_i(b_j) = \sum_{i,j} \hat{b}_j(a_i) = m \quad \square$$

We say that a COEA \mathcal{E} is *spectral* if \mathcal{E} is state-unique and for every $b \in \mathcal{E}$ there exists a context $\mathcal{A} = \{a_1, \dots, a_n\}$ such that

$$b = \lambda_1 a_1 \oplus \dots \oplus \lambda_n a_n$$

$\lambda_i \in [0, 1]$, $i = 1, \dots, n$. We denote the set of such $b \in \mathcal{E}$ corresponding to a fixed context \mathcal{A} by $\mathcal{E}(\mathcal{A})$. It can be shown that every COSEA is spectral [22]. A subset \mathcal{F} of a COEA \mathcal{E} is a *sub-COEA* if $0, 1 \in \mathcal{F}$, $a \in \mathcal{F}$ implies $d\lambda a \in \mathcal{F}$ for all $\lambda \in [0, 1]$ and if $a, b \in \mathcal{F}$ with $a \perp b$, then $a \oplus b \in \mathcal{F}$. A subset \mathcal{F} of a COSEA \mathcal{E} is a *sub-COSEA* if \mathcal{F} is a sub-COEA and if $a, b \in \mathcal{F}$ implies $a \circ b \in \mathcal{F}$. It is clear that if \mathcal{E} is a COEA (COSEA) then $\mathcal{E}(\mathcal{A})$ is a sub-COEA (sub-COSEA) for every $\mathcal{A} \in \mathcal{C}(\mathcal{E})$.

We close this section with some examples of COEA's and COSEA's. The first example comes from the quantum formalism. Let H be a complex Hilbert space and let $\mathcal{E}(H)$ be the set of operators on H satisfying $0 \leq A \leq I$ where we are using the usual operator order. For $A, B \in \mathcal{E}(H)$ we write $A \perp B$ if $A + B \leq I$ and in this case we define $A \oplus B = A + B$. For $\lambda \in [0, 1]$ and $A \in \mathcal{E}(H)$, $\lambda A \in \mathcal{E}(H)$ is the usual scalar multiplication for operators. It is easy to check that $(\mathcal{E}(H), 0, I, \oplus)$ is a COEA which we call a *Hilbertian COEA*. The sharp elements of $\mathcal{E}(H)$ are the projections on H . For $\phi \in H$ with $\phi \neq 0$, we denote the projection onto the one-dimensional subspace generated by ϕ as $P(\phi)$. Of course, $P(\phi) = P(\psi)$ if and only if $\phi = \alpha\psi$ for some $\alpha \in \mathbb{C}$, $\alpha \neq 0$. The elements of $S_1(\mathcal{E}(H))$ are precisely the $P(\phi)$, $\phi \in H$, $\phi \neq 0$ and $\mathcal{E}(H)$ is finite-dimensional if and only if H finite-dimensional. In this case, the contexts of $\mathcal{E}(H)$ correspond to the orthonormal bases of H so $\mathcal{C}(\mathcal{E}(H))$ is infinite if $\dim H \geq 2$. If $A \in S_1(\mathcal{E}(H))$ with $A = P(\phi)$ where $\|\phi\| = 1$, then \hat{A} is the unique state given by $\hat{A}(B) = \langle \phi, B\phi \rangle$ for all $B \in \mathcal{E}(H)$. Hence, $\mathcal{E}(H)$ is state-unique. It follows from the spectral theorem that $\mathcal{E}(H)$ is state-unique. Moreover, if $B = P(\psi)$, $\|\psi\| = 1$, then the transition probability becomes

$$\hat{A}(B) = \hat{B}(A) = |\langle \phi, \psi \rangle|^2$$

so $\mathcal{E}(H)$ is symmetric. If \mathcal{F} is a sub-COEA of $\mathcal{E}(H)$ for some H , we call \mathcal{F} a *sub-Hilbertian* COEA. An example of a sub-Hilbertian COEA is a von Neumann algebra of operators on H . These are also spectral and symmetric. For $A, B \in \mathcal{E}(H)$ define the product $A \circ B = A^{1/2}BA^{1/2}$ where $A^{1/2}$ is the unique positive square root of A . It is shown in [12, 13] that with the product $A \circ B$, $\mathcal{E}(H)$ becomes a COSEA. We also have that $A \circ B = B \circ A$ if and only if $AB = BA$ [14]; that is, A and B commute. We then call $\mathcal{E}(H)$ a *Hilbertian* COSEA, and any sub-COEA of $\mathcal{E}(H)$ is a *sub-Hilbertian* COSEA. As before, a von Neumann algebra on H is an example of a sub-Hilbertian COSEA.

Our next example comes from fuzzy probability theory [2, 8]. Let $\Omega, (\mathcal{A})$ be a measurable space in which singleton sets are measurable and let $\mathcal{E}(\Omega, \mathcal{A})$ be the set of measurable functions on Ω with values in $[0, 1] \subseteq \mathbb{R}$. If we define the operations $\oplus, \lambda f$ and $f \circ g = fg$ analogously as in the previous example, $\mathcal{E}(\Omega, \mathcal{A})$ becomes a COSEA. The elements of $\mathcal{E}(\Omega, \mathcal{A})$ are called fuzzy events and we call $\mathcal{E}(\Omega, \mathcal{A})$ a *classical* COSEA. The elements of $S(\mathcal{E}(\Omega, \mathcal{A}))$ are the characteristic functions (or equivalently, the sets in \mathcal{A}) and $S_1(\mathcal{E}(\Omega, \mathcal{A}))$ consists of the characteristic functions of the singleton sets (or equivalently, the elements of Ω). Notice that $\mathcal{E}(\Omega, \mathcal{A})$ is finite-dimensional if and only if Ω is finite and in this case there is only one context. Also, $\mathcal{E}(\Omega, \mathcal{A})$ is symmetric and spectral. Conversely, it is shown in [11] that if a finite-dimensional COEA (COSEA) \mathcal{E} has only one context, then \mathcal{E} is isomorphic to classical COEA (COSEA). We have seen that a classical COEA contains only one context while a quantum (Hilbertian) COEA possesses an infinite number of different contexts. Is there anything between? That is, can a finite-dimensional spectral COEA \mathcal{E} have a finite number, greater than one, of disjoint contexts [11]? The answer to this question is negative. In fact, if \mathcal{E} has more than one context, then it has uncountably many [17].

3 Commutants

In this section, \mathcal{E} will denote a finite-dimensional COSEA. For $\mathcal{F} \subseteq \mathcal{E}$, the *commutant* of \mathcal{F} is defined as

$$\mathcal{F}' = \{b \in \mathcal{E} : b \mid a \text{ for all } a \in \mathcal{F}\}$$

Notice that \mathcal{F}' is a sub-COSEA of \mathcal{E} . If $\mathcal{F} \subseteq \mathcal{G} \subseteq \mathcal{E}$ then $\mathcal{G}' \subseteq \mathcal{F}'$. We also have that $\mathcal{F} \subseteq \mathcal{F}''$, $\mathcal{F}' = \mathcal{F}'''$, $\mathcal{F}' \cap \mathcal{G}' \subseteq (\mathcal{F} \cap \mathcal{G})'$ and $(\mathcal{F} \cup \mathcal{G})' \subseteq \mathcal{F}' \cup \mathcal{G}'$ for all $\mathcal{F}, \mathcal{G} \subseteq \mathcal{E}$. We say that $\mathcal{F} \subseteq \mathcal{E}$ is *commutative* if $a \circ b = b \circ a$ for all $a, b \in \mathcal{F}$. Clearly, \mathcal{F} is commutative if and only if $\mathcal{F} \subseteq \mathcal{F}'$. It is shown in [11] that \mathcal{E} is commutative if and only if \mathcal{E} has only one context and hence is isomorphic to a classical COSEA. We call \mathcal{E}' the *center* of \mathcal{E} . Thus, $\mathcal{E} = \mathcal{E}'$ if and only if \mathcal{E} is isomorphic to a classical COSEA and \mathcal{E}' is a commutative sub-COSEA of \mathcal{E} . It is clear that $\{\lambda 1 : \lambda \in [0, 1]\} \subseteq \mathcal{E}'$. We say that \mathcal{E} is a *factor* if $\mathcal{E}' = \{\lambda 1 : \lambda \in [0, 1]\}$.

We now define the *direct sum* $\mathcal{E} = \mathcal{E}_1 \oplus \mathcal{E}_2$ of two COSEA's $(\mathcal{E}_1, 0_1, 1_1, \oplus)$, $(\mathcal{E}_2, 0_2, 1_2, \oplus)$. We define $(\mathcal{E}, 0, 1, \oplus)$ by

$$\mathcal{E} = \mathcal{E}_1 \times \mathcal{E}_2 = \{(a_1, a_2) : a_1 \in \mathcal{E}_1, a_2 \in \mathcal{E}_2\}$$

with $0 = (0_1, 0_2)$, $1 = (1_1, 1_2)$. If $a = (a_1, a_2)$, then $a' = (a'_1, a'_2)$. If $a = (a_1, a_2)$, $b = (b_1, b_2)$ then $a \perp b$ if $a_1 \perp b_1$, $a_2 \perp b_2$ and

$$a \oplus b = (a_1 \oplus b_1, a_2 \oplus b_2)$$

For $\lambda \in [0, 1]$ define $\lambda(a_1, a_2) = (\lambda a_1, \lambda a_2)$ and we define

$$(a_1, a_2) \circ (b_1, b_2) = (a_1 \circ b_1, a_2 \circ b_2)$$

It is easy to check that \mathcal{E} is a COSEA. We have that $(a_1, a_2) \leq (b_1, b_2)$ if and only if $a_1 \leq b_1, a_2 \leq b_2$ and

$$\mathcal{E}' = \{(a_1, a_2) : a_1 \in \mathcal{E}'_1, a_2 \in \mathcal{E}'_2\}$$

Clearly, $(a_1, a_2) \in S(\mathcal{E})$ if and only if $a_1 \in S(\mathcal{E}'_1)$ and $a_2 \in S(\mathcal{E}'_2)$.

Lemma 3.1. *Let $\mathcal{E} = \mathcal{E}'_1 \oplus \mathcal{E}'_2$. (i) $(a_1, a_2) \in S_1(\mathcal{E})$ if and only if $a_1 = 0_1$ and $a_2 \in S_1(\mathcal{E}'_2)$ or $a_2 = 0_2$ and $a_1 \in S_1(\mathcal{E}'_1)$. (ii) $\mathcal{A} \in \mathcal{C}(\mathcal{E})$ if and only if*

$$\mathcal{A} = \{(a_i, 0_2), (0_1, b_j)\}$$

where $\{a_i\} \in \mathcal{C}(\mathcal{E}'_1)$ and $\{b_j\} \in \mathcal{C}(\mathcal{E}'_2)$

Proof. (i) Necessity is clear. For sufficiency, suppose that $(a_1, a_2) \in S_1(\mathcal{E})$ and $a_1 \neq 0_1, a_2 \neq 0_2$. Then $(a_1, 0_2) \leq (a_1, a_2)$ but for $\lambda \in [0, 1]$ we have that

$$(a_1, 0_2) \neq (\lambda a_1, \lambda a_2) = \lambda(a_1, a_2)$$

which is a contradiction. Hence, $a_1 = 0_1$ or $a_2 = 0_2$. Clearly, if $a_1 \neq 0$, then $a_1 \in S_1(\mathcal{E}'_1)$ and if $a_2 \neq 0$, then $a_2 \in S_1(\mathcal{E}'_2)$. (ii) This follows from (i). \square

We shall need the following lemma to prove Theorem 3.3.

Lemma 3.2. (i) *If $a \mid c$ and $a \mid (c \oplus d)$ then $a \mid d$. (ii) If $c \leq b$ and $a \mid c, a \mid b$ then $a \mid (b \ominus c)$. (iii) If $c \leq b$ then $b \ominus c = (c \oplus b')'$. (iv) If \mathcal{F} is a sub-COSEA of \mathcal{E} and $b, c \in \mathcal{F}$ with $c \leq b$, then $b \ominus c \in \mathcal{F}$.*

Proof. (i) Let $b = c \oplus d$ so that $a \mid c$ and $a \mid b$. Now $c \oplus d \oplus b' = 1$ so $d = (c \oplus b')'$. Since $a \mid b, a \mid b'$ and since $a \mid c$ we have that $a \mid c \oplus b'$. Hence, $a \mid d$. (ii) Since $c \leq b$ we have that $b = c \oplus (b \ominus c)$. Since $a \mid c$ and $a \mid b$, by (i) $a \mid (b \ominus c)$. (iii) This follows from (i). (iv) Since $b, c \in \mathcal{F}$ we have that b' and $c \oplus b' \in \mathcal{F}$. Hence, by (iii).

$$b \ominus c = (c \oplus b')' \in \mathcal{F} \quad \square$$

Theorem 3.3. *A COSEA \mathcal{E} is isomorphic to a direct sum of two COSEA's if and only if there exists an $a \in S(\mathcal{E}) \cap \mathcal{E}'$ with $a \neq 0, 1$.*

Proof. If \mathcal{E} is isomorphic to a direct sum of two COSEA's, we can just as well assume that $\mathcal{E} = \mathcal{E}'_1 \oplus \mathcal{E}'_2$. We then have that $(1_1, 0_2) \in S(\mathcal{E}) \cap \mathcal{E}'$ and $(1_1, 0_2) \neq (1_1, 1_2) = 1$ and $(1_1, 0_2) \neq (0_1, 0_2) = 0$. Conversely, suppose $a \in S(\mathcal{E}) \cap \mathcal{E}'$ with $a \neq 0, 1$. Let

$$\mathcal{E}'_1 = \{a \circ b : b \in \mathcal{E}\}$$

and define $0_1 = a \circ 0 = 0$ and $1_1 = a \circ 1 = a$. For $a \circ b \in \mathcal{E}'_1$ define

$$(a \circ b)' = a \circ b' \in \mathcal{E}'_1$$

Define $a \circ b_1 \perp_1 a \circ b_2$ if $b_1 \perp b_2$ and in this case

$$a \circ b_1 \oplus_1 a \circ b_2 = a \circ (b_1 \oplus b_2) = a \circ b_1 \oplus a \circ b_2 \in \mathcal{E}'_1$$

It is easy to check that $(\mathcal{E}'_1, 0_1, 1_1, \oplus_1)$ is an effect algebra. Letting $\lambda(a \circ b) = a \circ (\lambda b)$ makes \mathcal{E}'_1 into a COSEA. Defining

$$(a \circ b) \circ_1 (a \circ b) = (a \circ b) \circ (a \circ c) = a \circ (b \circ c) \in \mathcal{E}'_1$$

we see that $a \circ b \mid_1 a \circ c$. We now show that $(\mathcal{E}_1, 0_1, 1_1, \oplus_1 \circ_1)$ is a COSEA. It is easy to verify that (S1) and (S2) hold. To verify (S3) suppose that $(a \circ b) \circ_1 (a \circ c) = 0$. Then

$$a \circ (b \circ c) = (a \circ b) \circ (a \circ c) = 0$$

Hence, $a \circ b \mid a \circ c$ so $a \circ b \mid_1 a \circ c$. To verify (S4) suppose that $a \circ b \mid_1 a \circ c$. Then $a \circ b \mid a \circ c$. Since $a = a \circ c \oplus a \circ c'$ and $a \circ b \mid a, a \circ b \mid a \circ c$ it follows from Lemma 3.2(i) that $a \circ b \mid a \circ c'$ so $a \circ b \mid (a \circ c)'$. Moreover, for all $d \in \mathcal{E}$ we have

$$(a \circ b) \circ [(a \circ c) \circ (a \circ d)] = [(a \circ b) \circ (a \circ c)] \circ (a \circ d)$$

The verification of (S5) and (S6) are straightforward. We conclude that \mathcal{E}_1 is a COSEA. Now $d' \in S(\mathcal{E})$ with $d' \neq 0, 1$ so letting $\mathcal{E}_1 = \{d' \circ b : b \in \mathcal{E}\}$ with similar definitions we have that $(\mathcal{E}_2, 0_2, 1_2, \oplus_2, \circ_2)$ is a COSEA. Every element of \mathcal{E} has the unique representation $b = a \circ b \oplus d' \circ b, a \circ b \in \mathcal{E}_1, d' \circ b \in \mathcal{E}_2$. Defining the map $J: \mathcal{E} \rightarrow \mathcal{E}_1 \oplus \mathcal{E}_2$ by $J(b) = (a \circ b, d' \circ b)$ it is straightforward to show that J is an isomorphism. \square

Since \mathcal{E} is spectral, every $b \in \mathcal{E}$ has a representation $b = \lambda_1 a_1 \oplus \cdots \oplus \lambda_n a_n$ for some $\{a_i\} \in \mathcal{C}(\mathcal{E})$, $\lambda_i \in [0, 1]$. We denote the set of effects that have such a representation relative to a context $\mathcal{A} \in \mathcal{C}(\mathcal{E})$ by $\mathcal{E}(\mathcal{A})$. It is clear that $\mathcal{E}(\mathcal{A})$ is a commutative sub-COSEA of \mathcal{E} . In fact, if b is as above and $c = \mu_1 a_1 \oplus \cdots \oplus \mu_n a_n, \mu_i \in [0, 1]$, then $b \perp c$ if and only if $\lambda_i + \mu_i \leq 1, i = 1, \dots, n$ and in this case

$$b \oplus c = (\lambda_1 + \mu_1) a_1 \oplus \cdots \oplus (\lambda_n + \mu_n) a_n$$

In general, we have

$$b \circ c = (\lambda_1 \mu_1) a_1 \oplus \cdots \oplus (\lambda_n \mu_n) a_n$$

In the representation for $b \in \mathcal{E}$, the λ_i need not be distinct but since the sum of sharp elements is sharp, we can write

$$b = \lambda'_1 c_1 \oplus \cdots \oplus \lambda'_m c_m \tag{3.1}$$

where $c_i \in S(\mathcal{E}), \lambda'_i \neq \lambda'_j, i \neq j$. The next result follows from Theorem 4.3 in [11].

Theorem 3.4. Any $b \in \mathcal{E}$ has a unique representation (3.1) where $\lambda'_i \in [0, 1], \lambda'_i \neq \lambda'_j, i \neq j, c_i \in S(\mathcal{E}), c_1 \oplus \cdots \oplus c_m = 1$ and $c_i \in \{b\}''$.

Theorem 3.5. In a COSEA \mathcal{E} , $a \mid b$ if and only if $a, b \in \mathcal{E}(\mathcal{A})$ for some $\mathcal{A} \in \mathcal{C}(\mathcal{E})$.

Proof. If $a, b \in \mathcal{E}(\mathcal{A})$, then clearly $a \mid b$. Conversely, suppose that $a \mid b$. By Theorem 3.4, we have $a = \oplus \lambda_i a_i, b = \oplus \mu_j b_j, \lambda_i \neq \lambda_j, \mu_i \neq \mu_j, i \neq j, a_i, b_i \in S(\mathcal{E})$ and $\oplus a_i = \oplus b_i = 1$. Moreover, by Theorem 3.4, $a_i \mid b_j$ for all i, j . Then $a_i \circ b_j \in S(\mathcal{E})$ and $\oplus a_i \circ b_j = 1$. Letting e_k be the nonzero $a_i \circ b_j$ we have that $e_k \in S(\mathcal{E})$ and $\oplus e_k = 1$. Then $a_i = \oplus \{e_k : e_k \leq a_i\}$ and similarly for the b_i . Reordering the λ_i and μ_i if necessary we can write $a = \oplus \lambda_i e_i, b = \oplus \mu_i e_i$. Finally, we can construct a context $\mathcal{A} = \{c_k\}$ such that $e_i = \oplus c_{k_i}$ for all i . Then $a = \oplus \lambda_i c_i, b = \oplus \mu_i c_i$ so that $a, b \in \mathcal{E}(\mathcal{A})$. \square

Lemma 3.6. If $a \in S_1(\mathcal{E}), b \in S(\mathcal{E})$, then $a \mid b$ if and only if $a \circ b = 0$ or $a \leq b$.

Proof. If $a \circ b = 0$ or $a \leq b$, then by Theorem 2.1, $a \mid b$. If $a \mid b$, then since $a \circ b \leq a$ we have that $a \circ b = \lambda a$ for some $\lambda \in [0, 1]$. Since $a \circ b \in S(\mathcal{E}), \lambda^2 a = \lambda a$ so $\lambda^2 = \lambda$. Hence, $\lambda = 0$ or $\lambda = 1$. If $\lambda = 0$, then $a \circ b = 0$. If $\lambda = 1$, then

$$a = a \circ b = b \circ a \leq b \tag{3.2}$$

Theorem 3.7. *If $a \in S_1(\mathcal{E})$, then*

$$\{a\}' = \left\{ b : b = \lambda a \oplus \bigoplus \lambda_i a_i, \{a, a_1, \dots, a_n\} \in \mathcal{C}(\mathcal{E}), \lambda, \lambda_i \in [0, 1] \right\} \quad (3.2)$$

Proof. If $b = \lambda a \oplus \bigoplus \lambda_i a_i$ as in (3.2), then clearly $b \mid a$. Conversely suppose $b \mid a$. By Theorem 3.4 we can write $b = \bigoplus \mu_i c_i$, $c_i \in S(\mathcal{E})$, $\mu_i \neq \mu_j$, $\mu_i \neq 0$, $c_i \circ c_j = 0$, $i \neq j$. Also by Theorem 3.4 we have that $a \mid c_i$ for all i so by Lemma 3.6 $a \circ c_i = 0$ or $a \leq c_i$. If $a \circ c_i = 0$ for all i , then form a context $\{a, a_1, \dots, a_n\}$ such that $b = 0a \oplus \bigoplus \lambda_i a_i$. Otherwise, there is a j such that $a \leq c_j$ and $a \circ c_i = 0$ for all $i \neq j$. We again form a context $\{a, a_1, \dots, a_n\}$ such that $b = \lambda a \oplus \bigoplus \lambda_i a_i$. \square

Theorem 3.8. *A COSEA \mathcal{E} is a factor if and only if \mathcal{E} is not isomorphic to the direct sum of two COSEA's.*

Proof. Suppose \mathcal{E} is a factor. If \mathcal{E} is isomorphic to a direct sum of COSEA's $\mathcal{E}_1, \mathcal{E}_2$, then by Theorem 3.3 there is an $a \in S(\mathcal{E}) \cap \mathcal{E}'$ with $a \neq 0, 1$. But then $a = \lambda 1$ for some $\lambda \in (0, 1)$. Since $a^2 = a$ we have that $\lambda^2 = \lambda$ so $\lambda = 0$ or $\lambda = 1$ which is a contradiction. Conversely, suppose \mathcal{E} is not a factor so that $\mathcal{E}' \neq \{\lambda 1 : \lambda \in [0, 1]\}$. Then there is a $b \in \mathcal{E}'$ with $b \neq \lambda 1$ for any $\lambda \in [0, 1]$. By Theorem 3.4, there exists an $a \in S(\mathcal{E}) \cap \{b\}''$ with $a \neq 0, 1$. Since $\{b\}' = \mathcal{E}$ we have that $a \in \{b\}'' = \mathcal{E}'$. By Theorem 3.3, \mathcal{E} is isomorphic to the direct sum of two COSEA's. \square

For $\mathcal{F} \subseteq \mathcal{E}$, if $a \in \mathcal{F} \cap S(\mathcal{E})$ with $a \neq 0$, we say that a is *minimal sharp* in \mathcal{F} if $b \in \mathcal{F} \cap S(\mathcal{E})$ and $b \leq a$, then $b = a$.

Theorem 3.9. *\mathcal{F} is a commutative sub-COSEA of \mathcal{E} if and only if there exist minimal sharp elements a_1, \dots, a_n in \mathcal{F} such that $a_1 \oplus \dots \oplus a_n = 1$ and*

$$\mathcal{F} = \{\lambda_1 a_1 \oplus \dots \oplus \lambda_n a_n : \lambda_i \in [0, 1]\} \quad (3.3)$$

Proof. If \mathcal{F} has the form (3.3), since $a_i \mid a_j$, $\mathcal{F} \subseteq \mathcal{F}'$ and it is easy to show that \mathcal{F} is a sub-COSEA. Conversely, suppose \mathcal{F} is a commutative sub-COSEA of \mathcal{E} . If $b \in \mathcal{F} \cap S(\mathcal{E})$ with $b \neq 0$ we show there exists a minimal sharp a in \mathcal{F} such that $a \leq b$. If b is minimal sharp in \mathcal{F} we are finished. Otherwise, there exists an $a_1 \in \mathcal{F} \cap S(\mathcal{E})$ with $a_n \neq 0$ and $a_1 < b$. If a_1 is minimal sharp in \mathcal{F} we are finished. Otherwise, there exists an $a_2 \in \mathcal{F} \cap S(\mathcal{E})$ with $a_2 \neq 0$ and $a_2 < a_1 < b$. This process must end because if $a_1 > a_2 > a_3 > \dots$ with, $a_i \in \mathcal{F} \cap S(\mathcal{E})$, $a_i \neq 0$, then letting $b_i = a_i \ominus a_{i+1}$, $i = 1, 2, \dots$, we have $b_i \in S(\mathcal{E})$ and $b_i \perp b_j$, $i \neq j$. Since \mathcal{E} is spectral, there exist $c_i \in S_1(\mathcal{E})$ such that $c_i \leq b_i$, $i = 1, 2, \dots$, but this contradicts the finite-dimensionality of \mathcal{E} . We conclude that for $b \in \mathcal{F} \cap S(\mathcal{E})$ with $b \neq 0$, there is a minimal sharp a in \mathcal{F} such that $a \leq b$. Let a_1, a_2, \dots, a_n be the minimal sharp elements of \mathcal{F} . Again, because of finite dimensionality there is a finite number of these. Moreover, we have $a_1 \oplus \dots \oplus a_n = 1$. If $d \in \mathcal{F}$, then Theorem 3.4 there exist $d_j \in \mathcal{F} \cap S(\mathcal{E})$ such that

$$d = \lambda_1 d_1 \oplus \dots \oplus \lambda_m d_m$$

where $\lambda_j \in [0, 1]$ and $d_1 \oplus \dots \oplus d_m = 1$. By our previous work $d_j = \bigoplus a_{i_j}$ so that $d = \mu_1 a_1 \oplus \dots \oplus \mu_n a_n$, $\mu_i \in [0, 1]$. \square

Corollary 3.10. *There exist minimal sharp elements a_1, \dots, a_n in \mathcal{E}' such that $a_1 \oplus \dots \oplus a_n = 1$ and*

$$\mathcal{E}' = \{\lambda_1 a_1 \oplus \dots \oplus \lambda_n a_n : \lambda_i \in [0, 1]\}$$

Lemma 3.11. *If a is a minimal sharp element of \mathcal{E}' and $\mathcal{F} = \{a \circ b : b \in \mathcal{E}\}$, then \mathcal{F} is a COSEA with unit a and \mathcal{F} is a factor.*

Proof. We have shown in the proof of Theorem 3.3 that \mathcal{F} is a COSEA with unit a . To show that \mathcal{F} is a factor, we must show that $\mathcal{F}' \cap \mathcal{F} = \{\lambda a : \lambda \in [0, 1]\}$. If $a \circ b \in \mathcal{F}' \cap \mathcal{F} \cap S(\mathcal{E})$, then $a \circ b \mid a \circ c$ for all $c \in \mathcal{E}$. We also have that $(a \circ b) \circ (a' \circ c) = 0$ so $a \circ b \mid a' \circ c$ for all $c \in \mathcal{E}$. Since $c = a \circ c \oplus a' \circ c$ we have $a \circ b \mid c$ so $a \circ b \in \mathcal{E}'$. Since $a \circ b \leq a$ and a is minimal sharp in \mathcal{E} we conclude that if $b \neq 0$ then $a \circ b = a$. Hence, the only sharp elements of $\mathcal{F}' \cap \mathcal{F}$ are 0 and a . Since every $c \in \mathcal{F}' \cap \mathcal{F}$ has the form $c = \lambda_1 c_1 \oplus \cdots \oplus \lambda_n c_n$, $\lambda_i \in [0, 1]$, $c_i \in S(\mathcal{F})$ we have that $c = \lambda a$, $\lambda \in [0, 1]$. Therefore, \mathcal{F} is a factor. \square

We can extend the definition of direct sum to more than two summands. We define

$$\mathcal{E}_1 \oplus \mathcal{E}_2 \oplus \mathcal{E}_3 = (\mathcal{E}_1 \oplus \mathcal{E}_2) \oplus \mathcal{E}_3$$

and of course, the placement of the parenthesis is immaterial. In a similar way, we define $\mathcal{E} = \mathcal{E}_1 \oplus \mathcal{E} \oplus \cdots \oplus \mathcal{E}_n$. For convenience, write $(a_1, \dots, a_n) \in \mathcal{E}$ as $a_1 \oplus \cdots \oplus a_n$, $a_i \in \mathcal{E}_i$, $i = 1, \dots, n$. We then have $a_i \circ a_j = 0$, $i \neq j$, and $1_1 \oplus \cdots \oplus 1_n = 1$. Also,

$$\mathcal{E}' = \{a_1 \oplus \cdots \oplus a_n : a_i \in \mathcal{E}'_i\}$$

Theorem 3.12. *Any finite-dimensional COSEA \mathcal{E} is isomorphic to the direct sum of a finite number of factors.*

Proof. By Corollary 3.10 there exist minimal sharp elements a_1, \dots, a_n in \mathcal{E}' with $a_1 \oplus \cdots \oplus a_n = 1$. By Lemma 3.11, $\mathcal{E}_i = \{a_i \circ b : b \in \mathcal{E}\}$ is a factor with unit a_i . Since every $b \in \mathcal{E}$ has the form

$$b = a_1 \circ b \oplus \cdots \oplus a_n \circ b$$

it follows that \mathcal{E} is isomorphic to $\mathcal{E}_1 \oplus \cdots \oplus \mathcal{E}_n$. \square

We close this section with a result about the state space of the direct sum. If V is a real vector space and $A_1, \dots, A_n \subseteq V$ we define the *convex hull* of a_1, \dots, a_n by

$$\begin{aligned} CH(A_1, \dots, A_n) \\ = \{\lambda_1 v_1 + \cdots + \lambda_n v_n : \lambda_i \geq 0, \sum \lambda_i = 1, v_i \in A_i, i = 1, \dots, n\} \end{aligned}$$

Theorem 3.13. $\Omega(\mathcal{E}_1 \oplus \cdots \oplus \mathcal{E}_n) = CH(\Omega(\mathcal{E}_1), \dots, \Omega(\mathcal{E}_n))$

Proof. We shall show that $\Omega(\mathcal{E}_1 \oplus \mathcal{E}_2) = CH(\Omega(\mathcal{E}_1), \Omega(\mathcal{E}_2))$ and the general result easily follows. If $\omega_1 \in \Omega(\mathcal{E}_1)$, $\omega_2 \in \Omega(\mathcal{E}_2)$, $\lambda \in [0, 1]$, $(a, b) \in \mathcal{E} = \mathcal{E}_1 \oplus \mathcal{E}_2$, define

$$\omega(a, b) = \lambda \omega_1(a) + (1 - \lambda) \omega_2(b)$$

To show that $\omega \in \Omega(\mathcal{E})$ we have that

$$\omega(1_1, 1_2) = \lambda \omega_1(1_1) + (1 - \lambda) \omega_2(1_2) = 1$$

and

$$\begin{aligned} \omega[(a_1, a_2) \oplus (b_1, b_2)] &= \omega[(a_1 \oplus b_1, a_2 \oplus b_2)] \\ &= \lambda \omega_1(a_1 \oplus b_1) + (1 - \lambda) \omega_2(a_2 \oplus b_2) \\ &= \lambda [\omega_1(a_1) + \omega_1(b_1)] + (1 - \lambda) [\omega_2(a_2) + \omega_2(b_2)] \end{aligned}$$

$$\begin{aligned}
&= [\lambda \omega_1(a_1) + (1 - \lambda) \omega_2(a_2)] + [\lambda \omega_1(b_1) + (1 - \lambda) \omega_2(b_2)] \\
&= \omega(a_1, a_2) + \omega(b_1, b_2)
\end{aligned}$$

Hence, $CH(\Omega(\mathcal{E}_1), \Omega(\mathcal{E}_2)) \subseteq \Omega(\mathcal{E}_1 \oplus \mathcal{E}_2)$. To show that $\Omega(\mathcal{E}_1 \oplus \mathcal{E}_2) \subseteq CH(\Omega(\mathcal{E}_1), \Omega(\mathcal{E}_2))$, let $\omega \in \Omega(\mathcal{E}_1 \oplus \mathcal{E}_2)$. If $\omega(1_1, 0) = 0$ then for $b \in \mathcal{E}_2$ define $\omega_2(b) = \omega(0_1, b)$. Since $\omega(0_1, 1_2) = 1$, $\omega_2 \in \Omega(\mathcal{E}_2)$ and we have that

$$\omega(a, b) = \omega((a, 0_2) \oplus (0_1, b)) = \omega(0_1, b) = \omega_2(b)$$

Similarly, if $\omega(0_1, 1_2) = 0$, then letting $\omega_1(a) = \omega(a, 0_2)$ we have that $\omega(a, b) = \omega_1(a)$. If $\omega(1_1, 0_2)$, $\omega(0_1, 1_2) \neq 0$, define $\omega_1 \in \Omega(\mathcal{E}_1)$, $\omega_2 \in \Omega(\mathcal{E}_2)$ by

$$\omega_1(a) = \frac{1}{\omega(1_1, 0_2)} \omega(a, 0_2), \quad \omega_2(b) = \frac{1}{\omega(0_1, 1_2)} \omega(0_1, b)$$

Then $\omega(1_1, 0_2) + \omega(0_1, 1_2) = \omega(1) = 1$ and

$$\omega(a, b) = \omega(a, 0_2) + \omega(0_1, b) = \omega(1_1, 0_2) \omega_1(a) + \omega(0_1, 1_2) \omega_2(b) \quad \square$$

4 Conditioning and Spectra

As before \mathcal{E} will denote a finite-dimensional COSEA and if $a \in S_1(\mathcal{E})$ then \hat{a} is the unique state on \mathcal{E} such that $\hat{a}(a) = 1$. If $b \in \mathcal{E}$ and $\omega \in \Omega(\mathcal{E})$ with $\omega(b) \neq 0$ we define the *conditional probability for ω given b* as $\omega(c | b) = \omega(b \circ c) / \omega(b)$ for every $c \in \mathcal{E}$. Notice that $\omega(\cdot | b)$ is indeed a state on \mathcal{E} .

Theorem 4.1. *Let $a \in S_1(\mathcal{E})$. (i) \hat{a} is the unique state on \mathcal{E} such that $a \circ b = \hat{a}(b)a$ for all $b \in \mathcal{E}$. (ii) \hat{a} is the unique state on \mathcal{E} such that $\hat{a}(b) = \hat{a}(a \circ b)$ for all $b \in \mathcal{E}$. (iii) If $\omega \in \Omega(\mathcal{E})$ with $\omega(a) \neq 0$, then $\omega(b | a) = \hat{a}(b)$ for all $b \in \mathcal{E}$.*

Proof. (i) Since $a \circ b \leq a$, there exists $\lambda_a(b) \in [0, 1]$ such that $a \circ b = \lambda_a(b)a$. Applying \hat{a} to both sides gives $\lambda_a(b) = \hat{a}(a \circ b)$. It is clear that $\lambda_a \in \Omega(\mathcal{E})$ and $\lambda_a(a) = 1$. Hence, $\lambda_a = \hat{a}$ so that $a \circ b = \hat{a}(b)a$ for all $b \in \mathcal{E}$. If $\omega \in \Omega(\mathcal{E})$ satisfies $a \circ b = \omega(b)a$ for all $b \in \mathcal{E}$, letting $b = a$ gives

$$a = a \circ a = \omega(a)a$$

Hence, $\omega(a) = 1$ so $\omega = \hat{a}$. Thus, \hat{a} is unique. (ii) By (i) we have that

$$\hat{a}(a \circ b) = \hat{a}(b)\hat{a}(a) = \hat{a}(b)$$

for all $b \in \mathcal{E}$. If $\omega \in \Omega(\mathcal{E})$ satisfies $\omega(b) = \omega(a \circ b)$ for all $b \in \mathcal{E}$, letting $b = 1$ gives $\omega(a) = \omega(1) = 1$ so that $\omega = \hat{a}$. (iii) If $\omega(a) \neq 0$, applying (i) gives

$$\omega(b | a) = \frac{\omega(a \circ b)}{\omega(a)} = \frac{\omega(\hat{a}(b)a)}{\omega(a)} = \hat{a}(b) \quad \square$$

From Theorem 4.1(iii) we have that $\hat{a}(b) = \omega(b | a)$ for all $\omega \in \Omega(\mathcal{E})$ with $\omega(a) \neq 0$. We conclude that \hat{a} is the *universal conditional probability given a* .

Let $\hat{\Omega}(\mathcal{E}) = \Omega(\mathcal{E}) \cup \{0\}$ where $0(b) = 0$ for all $b \in \mathcal{E}$. For all $a \in \mathcal{E}$ we define the *conditional probability map* $\gamma_a: \hat{\Omega}(\mathcal{E}) \rightarrow \hat{\Omega}(\mathcal{E})$ by $\gamma_a(0) = 0$ and for $\omega \neq 0$

$$\gamma_a(\omega) = \begin{cases} \omega(\cdot | a) & \text{if } \omega(a) \neq 0 \\ 0 & \text{if } \omega(a) = 0 \end{cases}$$

It is clear that $\gamma_0(\omega) = 0$ and $\gamma_1(\omega) = \omega$ for all $\omega \in \hat{\Omega}(\mathcal{E})$. The next result summarizes properties of γ .

Lemma 4.2. (i) If $a \in S_1(\mathcal{E})$, the \hat{a} is the unique nonzero fixed point of γ_a ; that is, $\gamma_a \omega = \omega$, $\omega \neq 0$ implies that $\omega = \hat{a}$. (ii) If $a \perp b$, $c \mid a$, $c \mid b$ then for all $\omega \in \Omega(\mathcal{E})$ we have that

$$\omega(a \oplus b) \gamma_{a \oplus b}(\omega)(c) = \omega(a) \gamma_a(\omega)(c) + \omega(b) \gamma_b(\omega)(c) \quad (4.1)$$

(iii) If $a \mid b$, then for all $\omega \in \widehat{\Omega}(\mathcal{E})$ we have that

$$\omega(a') \gamma_{a'}(\omega)(b) = \omega(b) - \omega(a) \gamma_a(\omega)(b) \quad (4.2)$$

(iv) For all $\omega \in \widehat{\Omega}(\mathcal{E})$ and $c \in \mathcal{E}$ we have that

$$\omega(a \circ b) \gamma_{a \circ b}(\omega)(c) = \omega[(a \circ b) \circ c] \quad (4.3)$$

and

$$\omega(a \circ b) [\gamma_b \gamma_a(\omega)](c) = \omega[a \circ (b \circ c)] \quad (4.4)$$

Proof. Conditions (4.1)–(4.4) clearly hold if $\omega = 0$. We thus assume that $\omega \in \Omega(\mathcal{E})$. (i) We have from Theorem 4.1(ii) that

$$\gamma_a(\hat{a})(b) = \hat{a}(b \mid a) = \hat{a}(a \circ b) = \hat{a}(b)$$

Hence, $\gamma_a(\hat{a}) = \hat{a}$. Now if $\gamma_a \omega = \omega$, then $\omega(a) \neq 0$ and for every $b \in \mathcal{E}$ we have that

$$\omega(b) = \gamma_a(\omega)(b) = \frac{\omega(a \circ b)}{\omega(a)}$$

We conclude that $\omega(a) = 1$ so that $\omega = \hat{a}$. (ii) If $\omega(a \oplus b) = 0$, then $\omega(a) = \omega(b) = 0$ so both sides of (4.1) are 0. If $\omega(a \oplus b) \neq 0$, then (4.1) is equivalent to

$$\begin{aligned} \omega[(a \oplus b) \circ c] &= \omega[c \circ (a \oplus b)] = \omega(c \circ a \oplus c \circ b) \\ &= \omega(c \circ a) + \omega(c \circ b) = \omega(a \circ c) + \omega(b \circ c) \end{aligned}$$

(iii) If $\omega(a') = 0$, then the left side of (4.2) is 0 and the right side is $\omega(b) = \omega(a \circ b)$. But $b = b \circ a \oplus b \circ a'$ and since $b \circ a' = a' \circ b \leq a'$ we have that $\omega(b \circ a') = 0$. Hence, $\omega(b) = \omega(a \circ b)$ so the right side is also 0. If $\omega(a') \neq 0$, then (4.2) is equivalent to

$$\begin{aligned} \omega(a' \circ b) &= \omega(b \circ a') = \omega(b) - \omega(b \circ a) = \omega(b) - \omega(a \circ b) \\ &= \omega(b) = \omega(a) \gamma_a(\omega)(b) \end{aligned}$$

(iv) If $\omega(a \circ b) = 0$, then both sides of (4.3) are 0. If $\omega(a \circ b) \neq 0$, then (4.3) follows directly. Since $b \circ c \leq b$, we have that $a \circ (b \circ c) \leq a \circ b$. Thus, if $\omega(a \circ b) = 0$ then both sides of (4.4) are 0. If $\omega(a \circ b) \neq 0$, then

$$\omega(a \circ b) [\gamma_b \gamma_a(\omega)](c) = \frac{\omega(a \circ b) \gamma_a(\omega)(b \circ c)}{\gamma_a(\omega)(b)} = \omega[a \circ (b \circ c)] \quad \square$$

If $\omega(a \oplus b) \neq 0$, then (4.1) shows that on $\{a, b\}'$ we have that $\gamma_{a \oplus b}$ is a convex combination

$$\gamma_{a \oplus b} = \frac{\omega(a)}{\omega(a) + \omega(b)} \gamma_a + \frac{\omega(b)}{\omega(a) + \omega(b)} \gamma_b$$

If $\omega(a') \neq 0$, then (4.2) implies that on $\{a\}'$ we have that

$$\gamma_{a'} = \frac{\omega - \omega(a)\gamma_a(\omega)}{1 - \gamma(a)}$$

If $a \mid b$, then (4.3) and (4.4) imply that

$$\gamma_b \gamma_a = \gamma_a \gamma_b = \gamma_{a \circ b}$$

We know that for $a \in S_1(\mathcal{E})$ there exists a unique $\omega \in \Omega(\mathcal{E})$ such that $\omega(a) = 1$. We now consider whether there are other effects with this property.

Theorem 4.3. *There exists a unique $\omega \in \Omega(\mathcal{E})$ for which $\omega(a) = 1$ if and only if there is a context $\{a_i\}$ such that*

$$a = a_1 \oplus \lambda_2 a_2 \oplus \cdots \oplus \lambda_n a_n \quad (4.5)$$

where $\lambda_i \in [0, 1]$.

Proof. If a has the form (4.5), then $\widehat{a}_1(a) = 1$. If $\omega \in \Omega(\mathcal{E})$ with $\omega(a) = 1$, then

$$\omega(a_1) + \sum_{i=2}^n \lambda_i \omega(a_i) = 1$$

If $\omega(a_j) \neq 0$ for some $j = 2, \dots, n$ then

$$1 = \omega(a_1) + \sum_{i=2}^n \lambda_i \omega(a_i) < \omega(a_1) + \sum_{i=2}^n \omega(a_i) = 1$$

which is a contradiction. Hence, $\omega(a_j) = 0$, $j = 2, \dots, n$. We conclude that $\omega(a_1) = 1$ so $\omega = \widehat{a}_1$ and \widehat{a}_1 is the unique state such that $\widehat{a}_1(a) = 1$. Conversely, suppose there exists a unique $\omega \in \Omega(\mathcal{E})$ such that $\omega(a) = 1$. Let $a = \oplus \lambda_i a_i$ for some $\{a_i\} \in \mathcal{C}(\mathcal{E})$, $\lambda_i \in [0, 1]$. Since $\omega(a) = 1$ we have that

$$\sum \lambda_i \omega(a_i) = \omega(a) = 1$$

If $\omega(a_j) \neq 0$ and $\lambda_j < 1$, then

$$1 = \sum \lambda_i \omega(a_i) < \sum \omega(a_i) = 1$$

which is a contradiction. Since $\omega(a_j) \neq 0$ for some j we have $\lambda_j = 1$ for some j . We can assume that $j = 1$ and write a in the form (4.5). We have that $\lambda_i < 1$, $i = 2, \dots, n$ because if $\lambda_i = 1$ then $\widehat{a}_i(a) = \widehat{a}_i(a) = 1$ which contradicts the uniqueness of ω . \square

Corollary 4.4. *If $a \in S(\mathcal{E})$, then there exists a unique $\omega \in \Omega(\mathcal{E})$ such that $\omega(a) = 1$ if and only if $a \in S_1(\mathcal{E})$.*

We say that $b \in \mathcal{E}$ is *dispersion-free* relative to $\omega \in \Omega(\mathcal{E})$ if $\omega(b^2) = \omega(b)^2$. Notice that if $b \in S(\mathcal{E})$, then $\omega(b^2) = \omega(b)^2$ if and only if $\omega(b) = 0$ or $\omega(b) = 1$. This terminology is due to the definition of dispersion as

$$\omega \left[(b - \omega(b)1)^2 \right] = \omega(b^2) - \omega(b)^2 \geq 0$$

We say that b is *constant almost everywhere* ω [a.e. (ω)] if $b = \lambda a \oplus c$, $\lambda \in [0, 1]$, where $a \in S(\mathcal{E})$, $a \circ c = 0$, $\omega(a) = 1$.

Theorem 4.5. An effect b is dispersion-free relative to $\omega \in \Omega(\mathcal{E})$ if and only if b is constant a.e. (ω).

Proof. If b is constant a.e. (ω), then $b = \lambda a \oplus c$, $a \in S(\mathcal{E})$, $a \circ c = 0$, $\omega(a) = 1$. Then $a \mid c$ and we have that $b^2 = \lambda^2 a \oplus c^2$. Since

$$a = a \circ c \oplus a \circ c' = a \circ c' = c' \circ a \leq c'$$

we have that $1 = \omega(a) \leq \omega(c')$. Hence, $\omega(c') = 1$ so that $\omega(c) = 0$. Since $c^2 \leq c$ and $\omega(c) = 0$ we conclude that $\omega(c^2) = 0$. Hence,

$$\omega(b^2) = \lambda^2 \omega(a) = \lambda^2 = \omega(b)^2$$

Conversely, suppose $\omega(b^2) = \omega(b)^2$. Let $b = \lambda_1 a_1 \oplus \cdots \oplus \lambda_n a_n$, $\lambda_i \in [0, 1]$, $\{a_i\} \in \mathcal{C}(\mathcal{E})$. Define the random variable $f(a_i) = \lambda_i$ with distribution $\omega(a_i)$. Then the expectation of f becomes

$$E_\omega(f) = \sum \lambda_i \omega(a_i) = \omega(b)$$

and

$$E_\omega(f^2) = \sum \lambda_i^2 \omega(a_i) = \omega(b^2) = \omega(b)^2 = E_\omega(f)^2$$

Hence,

$$E_\omega \left[(f - E_\omega(f))^2 \right] = E_\omega(f^2) - E(f)^2 = 0$$

Since $(f - E_\omega(f))^2 \geq 0$, $f = E_\omega(f)$ a.e. (ω). Therefore,

$$f(a_i) = E_\omega(f) = \omega(b) \text{ a.e. } (\omega)$$

We can assume that

$$f(a_1) = \cdots = f(a_m) = \omega(b)$$

and $\omega(a_{m+1}) = \cdots = \omega(a_n) = 0$. Letting $a = a_1 \oplus \cdots \oplus a_m$ and

$$c = \lambda_{m+1} a_{m+1} \oplus \cdots \oplus \lambda_n a_n$$

we have that $b = \omega(b)a \oplus c$ where $a \in S(\mathcal{E})$, $a \circ c = 0$, $\omega(a) = 1$. □

It follows from the proof of Theorem 4.5 that if a is constant a.e. (ω) then the constant is $\omega(a)$.

We say that $b \in \mathcal{E}$ has *eigeneffect* $a \in S_1(\mathcal{E})$ if $b \mid a$. Notice that $b \mid a$ if and only if $b \circ a = \hat{a}(b)a$. We call $\hat{a}(b)$ the *eigenvalue* corresponding to eigeneffect a . The set of eigeneffects for b is the *eigenspace* $S_1(b)$ and the set of eigenvalues for b is the *spectrum* $\sigma(b)$. Since \mathcal{E} is spectral, every $b \in \mathcal{E}$ can be written as $b = \lambda_1 a_1 \oplus \cdots \oplus \lambda_n a_n$, $\lambda_i \in [0, 1]$, $\{a_i\} \in \mathcal{C}(\mathcal{E})$. Since $b \mid a_i$, it follows that $a_i \in S_1(b)$ and $\lambda_i = \hat{a}_i(b) \in \sigma(b)$, $i = 1, \dots, n$. The different eigenvalues of b are unique but there may be various eigeneffects corresponding to the same eigenvalues. For example, if $\lambda_1 = \lambda_2$, then $a_1, a_2 \in S_1(\mathcal{E})$ correspond to λ_1 . More generally, in this case if $c \in S_1(\mathcal{E})$ and $c \leq a_1 \oplus a_2$ then c corresponds to λ_1 . It is also clear that if $a, b \in S_1(b)$ correspond to different eigenvalues, then $a \circ b = 0$. Moreover, $b \in S(\mathcal{E})$ if and only if $\sigma(b) \subseteq \{0, 1\}$ and $b \in S_1(\mathcal{E})$ if and only if $1 \in \sigma(b)$ and $S_1(b) = \{b\}$.

We define $m(b) = \min \{\lambda : \lambda \in \sigma(b)\}$ and $M(b) = \max \{\lambda : \lambda \in \sigma(b)\}$. Of course, $0 \leq m(b) \leq M(b) \leq 1$. We define the *numerical range* $r(b) = [m(b), M(b)]$ and the *norm* $\|b\| = M(b)$. It is clear that $\sigma(\lambda b) = \lambda \sigma(b)$, $r(\lambda b) = \lambda r(b)$ and $\|\lambda b\| = \lambda \|b\|$ for all $b \in \mathcal{E}$, $\lambda \in [0, 1]$.

Lemma 4.6. $r(b) = \{\omega(b) : \omega \in \Omega(\mathcal{E})\}$

Proof. Let $a_1, a_2 \in S_1(b)$ with $b \circ a_1 = m(b)a_1$ and $b \circ a_2 = M(b)a_2$. For $\lambda \in [0, 1]$ we define $\omega_\lambda \in \Omega(\mathcal{E})$ by $\omega_\lambda = \lambda \hat{a}_1 + (1 - \lambda) \hat{a}_2$. We then have

$$\begin{aligned} r(b) &= [m(b), M(b)] = \{\lambda m(b) + (1 - \lambda)M(b) : \lambda \in [0, 1]\} \\ &= \{\lambda \hat{a}_1(b) + (1 - \lambda) \hat{a}_2(b) : \lambda \in [0, 1]\} = \{\omega_\lambda(b) : \lambda \in [0, 1]\} \\ &\subseteq \{\omega(b) : \omega \in \Omega(\mathcal{E})\} \end{aligned}$$

Conversely, if $b = \lambda_1 a_1 \oplus \cdots \oplus \lambda_n a_n$, $\{a_i\} \in \mathcal{C}(\mathcal{E})$ then $\sigma(b) = \{\lambda_i\}$. If $\omega \in \Omega(\mathcal{E})$, then $\omega(b) = \sum \lambda_i \omega(a_i)$. Since $\sum \omega(a_i) = 1$ we have that

$$m(b) = \sum m(b) \omega(a_i) \leq \sum \lambda_i \omega(a_i) \leq \sum M(b) \omega(a_i) = M(b)$$

Hence $m(b) \leq \omega(b) \leq M(b)$ and we conclude that

$$\{\omega(b) : \omega \in \Omega(\mathcal{E})\} \subseteq r(b) \quad \square$$

Theorem 4.7. (i) $\|b\| = \max \{\omega(b) : \omega \in \Omega(\mathcal{E})\}$. (ii) If $b_1 \perp b_2$ then

$$\|b_1 \oplus b_2\| \leq \|b_1\| + \|b_2\|$$

(iii) $\|b\| = 0$ if and only if $b = 0$ (iv) If $a \leq b$ then $\|a\| \leq \|b\|$ and for all $a \in \mathcal{E}$, $a \leq \|a\| 1$. (v) $\|a \circ b\| \leq \|a\| \|b\|$.

Proof. (i) follows from Lemma 4.6. (ii) By (i) we have that

$$\begin{aligned} \|b_1 \oplus b_2\| &= \max \{\omega(b_1 \oplus b_2) : \omega \in \Omega(\mathcal{E})\} = \max \{\omega(b_1) + \omega(b_2) : \omega \in \Omega(\mathcal{E})\} \\ &\leq \max \{\omega(b_1) : \omega \in \Omega(\mathcal{E})\} + \max \{\omega(b_2) : \omega \in \Omega(\mathcal{E})\} \\ &= \|b_1\| + \|b_2\| \end{aligned}$$

(iii) We have that $b = 0$ if and only if $\sigma(b) = \{0\}$ which is equivalent to $\|b\| = 0$. (iv) If $a \leq b$, then there exists a $c \in \mathcal{E}$ such that $b = a \oplus c$. Hence, for all $\omega \in \Omega(\mathcal{E})$ we have that

$$\omega(a) \leq \omega(a) + \omega(c) = \omega(b)$$

It follows from (i) that $\|a\| \leq \|b\|$. Since $a = \lambda_1 a_1 \oplus \cdots \oplus \lambda_n a_n$, $\{a_i\} \in \mathcal{C}(\mathcal{E})$, $\sigma(a) = \{\lambda_i\}$ we have that

$$a = \lambda_1 a_2 \oplus \cdots \oplus \lambda_n a_n \leq M(a)(a_1 \oplus \cdots \oplus a_n) = M(a)1 = \|a\| 1$$

(v) By (iv) we have $b \leq \|b\| 1$ and hence, $a \circ b \leq \|b\| a$. Again by (iv) we conclude that

$$\|a \circ b\| \leq \| \|b\| a \| = \|a\| \|b\| \quad \square$$

5 Representation Theorems

Let \mathcal{E} be a finite-dimensional spectral COSEA. For $\mathcal{A} = \{a_i\} \in \mathcal{C}(\mathcal{E})$ define the complex linear space

$$\mathcal{H}(\mathcal{A}) = \left\{ \sum \alpha_i \widehat{a}_i : \alpha_i \in \mathbb{C} \right\}$$

For $x, y \in \mathcal{H}(\mathcal{A})$ with $x = \sum \alpha_i \widehat{a}_i$, $y = \sum \beta_i \widehat{a}_i$ define the inner product $\langle x, y \rangle = \sum \overline{\alpha}_i \beta_i$. Thus, $\mathcal{H}(\mathcal{A})$ is a complex Hilbert space that we call the *state space for context \mathcal{A}* . Of course, $\mathcal{H}(\mathcal{A})$ has orthonormal basis $\widehat{\mathcal{A}} = \{\widehat{a}_i : i = 1, \dots, n\}$ and $\dim \mathcal{H}(\mathcal{A}) = n$. The elements of $\widehat{\mathcal{A}}$ can be thought of as states in $\Omega(\mathcal{E})$ or as unit vectors in $\mathcal{H}(\mathcal{A})$ which again correspond to Hilbert space pure states. We now show that this dual role is consistent. For $b \in \mathcal{E}$ define the linear operator L_b on $\mathcal{H}(\mathcal{A})$ by $L_b = \sum \widehat{a}_j(b) P(\widehat{a}_j)$. Notice that L_b is a positive operator, $L_0 = 0$, $L_1 = I$, $L_{b'} = I - L_b$ and if $b \perp c$ then $L_{b \oplus c} = L_b + L_c$. We then have that

$$\langle \widehat{a}_i, L_b \widehat{a}_i \rangle = \langle \widehat{a}_i, \sum \widehat{a}_j(b) P(\widehat{a}_j) \widehat{a}_i \rangle = \langle \widehat{a}_i, \widehat{a}_i(b) \widehat{a}_i \rangle = \widehat{a}_i(b)$$

so the dual roles are consistent. It is easy to see that $L: \mathcal{E} \rightarrow \mathcal{E}(\mathcal{H}(\mathcal{A}))$ need not be injective or surjective and does not preserve sharpness. Moreover, all the L_b , $b \in \mathcal{E}$, commute so they do not convey quantum interference. One can say that L gives a distorted partial view of \mathcal{E} . The reason for this is that we are only employing a single context \mathcal{A} . Each context gives a partial view and in order to obtain a total view, they must all be considered.

In order to consider several contexts together, we introduce a method to compare them. We say that \mathcal{E} is *comparable* if for every $\mathcal{A}, \mathcal{B} \in \mathcal{C}(\mathcal{E})$ there exists a unitary operator $U_{\mathcal{A}\mathcal{B}}: \mathcal{H}(\mathcal{A}) \rightarrow \mathcal{H}(\mathcal{B})$ such that

$$\left| \langle U_{\mathcal{A}\mathcal{B}} \widehat{a}, \widehat{b} \rangle \right|^2 = \widehat{a}(b) \quad (5.1)$$

for all $a \in \mathcal{A}$, $b \in \mathcal{B}$ and

$$U_{\mathcal{B}\mathcal{C}} U_{\mathcal{A}\mathcal{B}} = U_{\mathcal{A}\mathcal{C}} \quad (5.2)$$

for all $\mathcal{C} \in \mathcal{C}(\mathcal{E})$. Notice that if \mathcal{E} is comparable, then any two contexts in \mathcal{E} have the same cardinality.

We now justify why we assume that $\mathcal{H}(\mathcal{A})$ is a complex Hilbert space instead of a real space which may seem to be more natural. In many situations, there is an underlying symmetry group that we would like to represent on \mathcal{E} . This is most accurately accomplished by employing a unitary representation of the group on $\mathcal{H}(\mathcal{A})$ for some $\mathcal{A} \in \mathcal{C}(\mathcal{E})$. For a unitary representation, we need $\mathcal{H}(\mathcal{A})$ to be complex. Moreover, it is desirable for the representation to be context independent. This motivates requiring that \mathcal{E} is comparable because in this case the representations for different contexts are unitarily equivalent.

Lemma 5.1. *If \mathcal{E} is comparable, then (i) $U_{\mathcal{A}\mathcal{A}} = I$, (ii) $U_{\mathcal{A}\mathcal{B}} = U_{\mathcal{B}\mathcal{A}}^*$, (iii) $|\langle U_{\mathcal{A}\mathcal{B}} \widehat{a}, U_{\mathcal{C}\mathcal{B}} \widehat{c} \rangle|^2 = \widehat{a}(c)$.*

Proof. (i) Applying (5.2) gives $U_{\mathcal{A}\mathcal{A}} U_{\mathcal{A}\mathcal{A}} = U_{\mathcal{A}\mathcal{A}}$. Multiplying by $U_{\mathcal{A}\mathcal{A}}^*$ gives $U_{\mathcal{A}\mathcal{A}} = I$. (ii) By (5.2) we have that

$$U_{\mathcal{B}\mathcal{A}} U_{\mathcal{A}\mathcal{B}} = U_{\mathcal{A}\mathcal{A}} = I$$

Hence, $U_{\mathcal{A}\mathcal{B}} = U_{\mathcal{B}\mathcal{A}}^*$. (iii) Applying (5.1), (5.2) and (ii) gives

$$\begin{aligned} |\langle U_{\mathcal{A}\mathcal{B}} \widehat{a}, U_{\mathcal{C}\mathcal{B}} \widehat{c} \rangle|^2 &= |\langle U_{\mathcal{C}\mathcal{B}}^* U_{\mathcal{A}\mathcal{B}} \widehat{a}, \widehat{c} \rangle|^2 = |\langle U_{\mathcal{B}\mathcal{C}} U_{\mathcal{A}\mathcal{B}} \widehat{a}, \widehat{c} \rangle|^2 \\ &= |\langle U_{\mathcal{A}\mathcal{C}} \widehat{a}, \widehat{c} \rangle|^2 = \widehat{a}(c) \end{aligned} \quad \square$$

For $b \in \mathcal{E}(\mathcal{B})$ with $b = \lambda_1 b_1 \oplus \cdots \oplus \lambda_n b_n$, define $\tilde{b} \in \mathcal{H}(\mathcal{B})$ by $\tilde{b} = \sum \lambda_i P(\hat{b}_i)$. For comparable \mathcal{E} define $\tilde{U}_{\mathcal{B}\mathcal{A}}: \mathcal{E}(\mathcal{H}(\mathcal{B})) \rightarrow \mathcal{E}(\mathcal{H}(\mathcal{A}))$ by

$$\tilde{U}_{\mathcal{B}\mathcal{A}} = U_{\mathcal{B}\mathcal{A}} B U_{\mathcal{A}\mathcal{B}} = U_{\mathcal{B}\mathcal{A}} B U_{\mathcal{B}\mathcal{A}}^*$$

We say that \mathcal{E} is *strongly comparable* if \mathcal{E} is comparable and if $b_1 \perp b_2$ with $b_1 \in \mathcal{E}(\mathcal{A})$, $b_2 \in \mathcal{E}(\mathcal{B})$, $b_1 \oplus b_2 \in \mathcal{E}(\mathcal{C})$, then

$$(b_1 \oplus b_2)^\sim = \tilde{U}_{\mathcal{A}\mathcal{C}} \tilde{b}_1 \oplus \tilde{U}_{\mathcal{B}\mathcal{C}} \tilde{b}_2 \quad (5.3)$$

We see that (5.3) is a reasonable requirement which postulates that \oplus is independent of its Hilbert space representation.

Theorem 5.2. *A finite-dimensional COEA \mathcal{E} is isomorphic to a finite-dimensional Hilbertian sub-COEA if and only if \mathcal{E} is spectral and strongly comparable.*

Proof. Suppose \mathcal{E} is isomorphic to a Hilbertian sub-COEA \mathcal{F} . For simplicity we can assume that $\mathcal{E} = \mathcal{F}$. It is clear that \mathcal{F} is state-unique. By the spectral theorem, if $b \in \mathcal{F}$, then $b = \sum \lambda_i a_i$ where $a_i \in S_1(\mathcal{E}(H))$ are polynomial functions of b . Hence, $a_i \in \mathcal{F}$ so \mathcal{E} is spectral. To show that \mathcal{E} is comparable, let $\mathcal{A} = \{a_i\}$, $\mathcal{B} = \{b_i\}$ be contexts in \mathcal{E} . Then $\{\hat{a}_i\}$, $\{\hat{b}_i\}$ are orthonormal bases of \mathcal{H} . Define $U_{\mathcal{A}\mathcal{B}}: \mathcal{H}(\mathcal{A}) \rightarrow \mathcal{H}(\mathcal{B})$ by

$$U_{\mathcal{A}\mathcal{B}} \hat{a}_i = \sum_j \langle \hat{b}_j, \hat{a}_i \rangle \hat{b}_j$$

and extend by linearity. It is clear that $U_{\mathcal{A}\mathcal{B}}$ is unitary. Also, (5.1) holds because

$$\left| \langle U_{\mathcal{A}\mathcal{B}} \hat{a}_i, \hat{b}_j \rangle \right|^2 = \left| \langle \hat{b}_j, \hat{a}_i \rangle \right|^2 = \hat{a}_i(b_j)$$

If $\mathcal{C} = \{c_i\}$ is another context, we have that

$$\begin{aligned} U_{\mathcal{B}\mathcal{C}} U_{\mathcal{A}\mathcal{B}} \hat{a}_i &= \sum_j \langle \hat{b}_j, \hat{a}_i \rangle U_{\mathcal{B}\mathcal{C}} \hat{b}_j = \sum_{j,k} \langle \hat{b}_j, \hat{a}_i \rangle \langle \hat{c}_k, \hat{b}_j \rangle \hat{c}_k \\ &= \sum_k \langle \hat{c}_k, \hat{a}_i \rangle \hat{c}_k = U_{\mathcal{A}\mathcal{C}} \hat{a}_i \end{aligned}$$

Hence, (5.2) holds so \mathcal{E} is comparable. In this case, if $a \in \mathcal{E}$ then $a = \tilde{a}$ and $\tilde{U}_{\mathcal{A}\mathcal{C}} = \tilde{U}_{\mathcal{B}\mathcal{C}} = I$ so clearly \mathcal{E} is strongly comparable.

Conversely, suppose \mathcal{E} is spectral and strongly comparable. Fix $\mathcal{A} \in \mathcal{C}(\mathcal{E})$ and let $J: \mathcal{E} \rightarrow \mathcal{E}(\mathcal{H}(\mathcal{A}))$ be defined by

$$J(b) = \tilde{U}_{\mathcal{B}\mathcal{A}}(\tilde{b})$$

where $b \in \mathcal{E}(\mathcal{B})$, $\mathcal{B} \in \mathcal{C}(\mathcal{E})$. We first show that $J(b)$ is well-defined. That is, we need to show $J(b)$ is independent of the context \mathcal{B} containing b . Suppose $b \in \mathcal{E}(\mathcal{B}) \cap \mathcal{E}(\mathcal{C})$. Letting $b_1 = 0$, $b_2 = b$, we have that $b_1 \in \mathcal{E}(\mathcal{B})$, $b_2 \in \mathcal{E}(\mathcal{C})$ and $b_1 \oplus b_2 = b \in \mathcal{E}(\mathcal{C})$. By (5.3) we have that

$$\tilde{b} = (b_1 \oplus b_2)^\sim = \tilde{U}_{\mathcal{B}\mathcal{C}}(\tilde{b}_1) \oplus \tilde{U}_{\mathcal{B}\mathcal{C}}(\tilde{b}_2) = \tilde{U}_{\mathcal{B}\mathcal{C}}(\tilde{b})$$

Therefore

$$\tilde{U}_{\mathcal{C}\mathcal{A}}(\tilde{b}) = \tilde{U}_{\mathcal{C}\mathcal{A}} \tilde{U}_{\mathcal{B}\mathcal{C}}(\tilde{b}) = \tilde{U}_{\mathcal{B}\mathcal{A}}(\tilde{b})$$

Hence, $J(b)$ is well-defined. We now show that J is injective. Let $b \in \mathcal{E}(\mathcal{B})$ with $b = \lambda_1 b_1 \oplus \cdots \oplus \lambda_n b_n$, $c \in \mathcal{E}(\mathcal{C})$ with $c = \mu_1 c_1 \oplus \cdots \oplus \mu_n c_n$ and suppose that $J(b) = J(c)$. Then $\tilde{U}_{\mathcal{B}\mathcal{A}}(\tilde{b}) = \tilde{U}_{\mathcal{C}\mathcal{A}}(\tilde{c})$ or equivalently

$$U_{\mathcal{B}\mathcal{A}}\tilde{b}U_{\mathcal{A}\mathcal{B}} = U_{\mathcal{C}\mathcal{A}}\tilde{c}U_{\mathcal{A}\mathcal{C}}$$

This implies that

$$\tilde{b} = U_{\mathcal{A}\mathcal{B}}U_{\mathcal{C}\mathcal{A}}\tilde{c}U_{\mathcal{A}\mathcal{C}}U_{\mathcal{B}\mathcal{A}} = U_{\mathcal{C}\mathcal{B}}\tilde{c}U_{\mathcal{B}\mathcal{C}}$$

which gives $U_{\mathcal{B}\mathcal{C}}\tilde{b} = \tilde{c}U_{\mathcal{B}\mathcal{C}}$. We conclude that

$$\tilde{c}(U_{\mathcal{B}\mathcal{C}}\hat{b}_i) = (U_{\mathcal{B}\mathcal{C}}\tilde{b})\hat{b}_i = \lambda_i U_{\mathcal{B}\mathcal{C}}\hat{b}_i$$

Hence, $U_{\mathcal{B}\mathcal{C}}\hat{b}_i$ an eigeneffect of \tilde{c} with corresponding eigenvalue λ_i . But the eigenvalues of \tilde{c} are μ_j with corresponding eigeneffects \hat{c}_j . Therefore, $\lambda_i = \mu_j$ for some j and $U_{\mathcal{B}\mathcal{C}}\hat{b}_i = \hat{c}_j$. Since

$$\hat{b}_i(c_k) = \left| \left\langle U_{\mathcal{B}\mathcal{C}}\hat{b}_i, \hat{c}_k \right\rangle \right|^2 = \left| \langle \hat{c}_j, \hat{c}_k \rangle \right|^2 = \delta_{ij}$$

we have that $\hat{b}_i(c_j) = 1$. We conclude that $b_j = c_i$ for all i so $b = c$. We now show that $J(b_1 \oplus b_2) = J(b_1) \oplus J(b_2)$. Suppose that $b_1 \perp b_2$ with $b_1 \in \mathcal{E}(\mathcal{B})$, $b_2 \in \mathcal{E}(\mathcal{D})$, $b_1 \oplus b_2 \in \mathcal{E}(\mathcal{C})$. By strong comparability we have that

$$(b_1 \oplus b_2)^\sim = \tilde{U}_{\mathcal{B}\mathcal{C}}(\tilde{b}_1) \oplus \tilde{U}_{\mathcal{D}\mathcal{C}}(\tilde{b}_2)$$

Hence,

$$\begin{aligned} J(b_1 \oplus b_2) &= \tilde{U}_{\mathcal{C}\mathcal{A}}[(b_1 \oplus b_2)^\sim] = \tilde{U}_{\mathcal{C}\mathcal{A}}[\tilde{U}_{\mathcal{B}\mathcal{C}}(\tilde{b}_1) \oplus \tilde{U}_{\mathcal{D}\mathcal{C}}(\tilde{b}_2)] \\ &= \tilde{U}_{\mathcal{C}\mathcal{A}}\tilde{U}_{\mathcal{B}\mathcal{C}}(\tilde{b}_1) \oplus \tilde{U}_{\mathcal{C}\mathcal{A}}\tilde{U}_{\mathcal{D}\mathcal{C}}(\tilde{b}_2) = \tilde{U}_{\mathcal{B}\mathcal{A}}(\tilde{b}_1) \oplus \tilde{U}_{\mathcal{D}\mathcal{A}}(\tilde{b}_2) \\ &= J(b_1) \oplus J(b_2) \end{aligned}$$

If $\lambda \in [0, 1]$, $b \in \mathcal{E}(\mathcal{B})$, then

$$J(\lambda b) = \tilde{U}_{\mathcal{B}\mathcal{A}}((\lambda b)^\sim) = \tilde{U}_{\mathcal{B}\mathcal{A}}(\lambda \tilde{b}) = \lambda \tilde{U}_{\mathcal{B}\mathcal{A}}(\tilde{b}) = \lambda J(b)$$

It is easy to check that the range of J is a sub-COEA of $\mathcal{E}(\mathcal{H}(\mathcal{A}))$. □

We now consider representations of a finite-dimensional COSEA \mathcal{E} . We first need some preliminary lemmas. We saw in Theorem 3.4 that any $a \in \mathcal{E}$ with $a \neq 0$ has a unique representation $a = \lambda_1 c_1 \oplus \cdots \oplus \lambda_n c_n$, $\lambda_i \neq 0$, $\lambda_i \neq \lambda_j$, $i \neq j$, and $c_i \in S(\mathcal{E})$. We denote by $\lceil a \rceil$ the smallest sharp element that dominates a .

Lemma 5.3. $\lceil a \rceil$ exists and $\lceil a \rceil = c_1 \oplus \cdots \oplus c_n$.

Proof. Let $c = c_1 \oplus \cdots \oplus c_n$. Then $c \in S(\mathcal{E})$ and $a \leq c$. Suppose $b \in S(\mathcal{E})$ and $a \leq b$. Then $a \circ b = b \circ a = a$. Hence, $b \mid c_i$ and

$$\lambda_1 c_1 \oplus \cdots \oplus \lambda_n c_n = a = a \circ b = \lambda_1 c_1 \circ b \oplus \cdots \oplus \lambda_n c_n \circ b \tag{5.4}$$

Now $c_i \circ b \leq c_i$ and if $c_i \circ b < c_i$ we would contradict (5.4). Hence, $c_i \circ b = c_i$ so that

$$c \circ b = \oplus(c_i \circ b) = \oplus c_i = c$$

It follows that $c \leq b$ so that $c = \lceil a \rceil$. □

We say that $a \in \mathcal{E}$ is *pseudo-invertible* if there exists a $b \in \mathcal{E}$ such that $\lceil b \rceil = \lceil a \rceil$, $\|b\| = 1$ and

$$a \circ b = b \circ a = \lambda \lceil a \rceil$$

for some $\lambda \in [0, 1]$. We then call b a *pseudo-inverse* for a . (A slightly different definition as well as a version of the next lemma are given in [21].) We denote the smallest nonzero eigenvalue of a by $\lambda(a)$.

Lemma 5.4. *If $a \neq 0$, then a has a unique pseudo-inverse and $\lambda = \lambda(a)$.*

Proof. If $a \neq 0$, as before a has the unique representation $a = \lambda_1 c_1 \oplus \cdots \oplus \lambda_n c_n$, $\lambda_i \neq 0$, $\lambda_i \neq \lambda_j$, $i \neq j$, $c_i \in S(\mathcal{E})$. Letting

$$b = \lambda(a) \left(\frac{1}{\lambda_1} c_1 \oplus \cdots \oplus \frac{1}{\lambda_n} c_n \right)$$

we have from Lemma 5.3 that

$$a \circ b = b \circ a = \lambda(a)(c_1 \oplus \cdots \oplus c_n) = \lambda(a) \lceil a \rceil$$

Moreover, $\|b\| = 1$, $\lceil b \rceil = \lceil a \rceil = c_1 \oplus \cdots \oplus c_n$. For uniqueness, suppose $\lceil d \rceil = \lceil a \rceil$, $\|d\| = 1$ and $a \circ d = d \circ a = \lambda \lceil a \rceil$. Then $d = \mu_1 c_1 \oplus \cdots \oplus \mu_n c_n$ and

$$\mu_1 \lambda_1 c_1 \oplus \cdots \oplus \mu_n \lambda_n c_n = a \circ d = \lambda \lceil a \rceil$$

This implies that $\mu_i \lambda_i = \lambda$ for all i . Hence, $\mu_i = \lambda / \lambda_i$. Since $\|d\| = 1$ we have $M(d) = 1$ which implies that

$$\frac{\lambda}{\lambda(a)} = \frac{\lambda}{\min(\lambda_i)} = \max \left(\frac{\lambda}{\lambda_i} \right) = \max(\mu_i) = \|d\| = 1$$

Therefore, $\lambda(a) = \lambda$ and $\mu_i = \lambda(a) / \lambda_i$ so $d = b$. □

We denote the unique pseudo-inverse of a by a^{-1} . If $a \neq 0$, $\mu > 0$ and $\mu a \in \mathcal{E}$, then it is easy to show that $(\mu a)^{-1} = a^{-1}$. It follows that $(a^{-1})^{-1} = a / \|a\|$ and $((a^{-1})^{-1})^{-1} = a^{-1}$. We can interpret a^{-1} operationally as the effect that reverses a without interference but with a reduction of intensity by a factor $\lambda(a)$. If $\lceil a \rceil = 1$, we say that a is *invertible* and a^{-1} is the *inverse* of a . We say that \mathcal{E} is *inverse-preserving* if whenever a and b are invertible, then $a \circ b$ is as well and $(a \circ b)^{-1} = a^{-1} \circ b^{-1}$. Notice that the order of a^{-1} and b^{-1} on the right is a bit unexpected but this is the correct order for a sequential product $a \circ b$ in which a is measured first. It is clear that a classical COSEA is inverse-preserving. That a Hilbertian sub-COSEA is also will be shown in Theorem 5.6.

Lemma 5.5. (i) $a \in \mathcal{E}$ is invertible if and only if a does not have a zero eigenvalue. (ii) If $a \perp b$ and a is invertible then $a \oplus b$ is invertible.

Proof. (i) If $0 \in \sigma(a)$, then $\lceil a \rceil \neq 1$ so a is not invertible. If $0 \notin \sigma(a)$, then $\lceil a \rceil = 1$ so a is invertible. (ii) If a is invertible, the $\lceil a \rceil = 1$. Suppose $a \oplus b$ is not invertible. Then $\lceil a \oplus b \rceil \neq 1$ so there exists a $c \in S_1(\mathcal{E})$ such that

$$c \circ a \oplus c \circ b = c \circ (a \oplus b) = (a \oplus b) \circ c = 0$$

Hence, $c \circ a = 0$ which contradicts $\lceil a \rceil = 1$. □

When we consider a sub-Hilbertian COSEA $\mathcal{F} \subseteq \mathcal{E}(H)$ we are assuming the standard sequential product $A \circ B = A^{1/2} B A^{1/2}$ on \mathcal{F} .

Theorem 5.6. *A finite-dimensional COSEA \mathcal{E} is isomorphic to a finite-dimensional sub-Hilbertian COSEA $\mathcal{F} \subseteq \mathcal{E}(H)$ if and only if \mathcal{E} is strongly comparable and inverse-preserving.*

Proof. Suppose \mathcal{E} is COSEA isomorphic to $\mathcal{F} \subseteq \mathcal{E}(H)$. For simplicity, we can assume that $\mathcal{E} = \mathcal{F}$. We have shown in Theorem 5.2 that \mathcal{E} is strongly comparable. To show that \mathcal{E} is inverse-preserving, suppose that $A, B \in \mathcal{E}$ are invertible. It follows from Lemma 5.5(i) that A and B are invertible in the usual operator sense. To avoid confusion, denote the usual operator inverse of A by \widehat{A} . We then have that

$$(A \circ B)^\wedge = (A^{1/2} B A^{1/2})^\wedge = \widehat{A}^{1/2} \widehat{B} \widehat{A}^{1/2} = \widehat{A} \circ \widehat{B} \quad (5.5)$$

Writing A^{-1} as we previously define it we have that

$$A \circ A^{-1} = A^{-1} \circ A = \lambda(A)I$$

Therefore, $A^{-1} = \lambda(A)\widehat{A}$. Hence, $\lambda(A)\widehat{A} \in \mathcal{E}$ although $\widehat{A} \notin \mathcal{E}$ in general. Similarly, $B^{-1} = \lambda(B)\widehat{B} \in \mathcal{E}$ and we can rewrite (5.5) as

$$A^{-1} \circ B^{-1} = \lambda(A)\lambda(B)\widehat{A} \circ \widehat{B} = \lambda(A)\lambda(B)(A \circ B)^\wedge = (A \circ B)^{-1}$$

Hence, $(A \circ B)^{-1}$ exists and equals $\lambda(A)\lambda(B)(A \circ B)^\wedge$.

Conversely, suppose \mathcal{E} is strongly comparable and inverse-preserving. We have previously observed that \mathcal{E} is automatically spectral. Applying Theorem 5.2 there exists a COSEA isomorphism J from \mathcal{E} onto a Hilbertian sub-COSEA \mathcal{F} of $\mathcal{E}(H)$. Define the product $J(a) \cdot J(b) = J(a \circ b)$ on \mathcal{F} . It is shown in [11] that \mathcal{F} becomes a COSEA under this product. If $a \in \mathcal{E}$ is invertible, then $J(a)$ is invertible with $J(a)^{-1} = J(a^{-1})$. Indeed, $\|J(a^{-1})\| = 1$, $J(a^{-1}) \mid J(a)$ and

$$J(a) \cdot J(a^{-1}) = J(a \circ a^{-1}) = J(\lambda(a)1) = \lambda(a)I$$

If \mathcal{E} is inverse preserving, then \cdot is also inverse preserving because if $J(a)$ and $J(b)$ are invertible, then a and b are invertible and

$$\begin{aligned} [J(a) \cdot J(b)]^{-1} &= [J(a \circ b)]^{-1} = J[(a \circ b)^{-1}] = J(a^{-1} \circ b^{-1}) \\ &= J(a^{-1}) \cdot J(b^{-1}) = J(a)^{-1} \cdot J(b)^{-1} \end{aligned}$$

We conclude that $\mathcal{F} \subseteq \mathcal{E}(H)$ is an inverse preserving COSEA with sequence product. But \mathcal{F} is also an inverse preserving COSEA under the standard sequential product \circ . It follows from Theorem 5.19 in [21] that $J(a) \cdot J(b) = J(a) \circ J(b)$. Hence, $J: \mathcal{E} \rightarrow \mathcal{F}$ is a COSEA isomorphism. \square

6 Closing Comments

A natural question the reader may ask is: ‘‘What is the relationship between contexts as discussed here and the concept of contextuality considered in the literature [1, 19, 20]?’’ We shall devote a few sentences to this question and leave a more complete investigation to a future work. The notion of contextuality is based on an ontological model for a quantum system. Such a model is described by a measurable space (Λ, Σ) where Λ is the set of pure states for the system. Preparation procedures, state transformations and measurements are defined by stochastic maps on Λ that satisfy certain conditions. One of the main assumptions is that these maps combine to reproduce the experimental statistics of the system in terms of conditional probabilities. We define preparation, transformation and measurement non-contextuality

when these stochastic maps satisfy injectiveness properties. Our point is that the concept of contexts can be employed to construct such ontological models by defining the stochastic maps on contexts. Conversely, the stochastic maps for an ontological model will have their supports precisely on the contexts that we have defined in this paper.

Finally, we should mention that other approaches to the mathematical foundations of quantum mechanics have been recently explored. In particular, there have been recent efforts to provide a new foundation for the Hilbert space framework of quantum theory [3, 4, 16]. The main difference is that these works emphasize the role of composite systems and general transformations, while the COSEA formalism focuses on individual systems and on transformations induced by conditioning with sharp effects.

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